



Stückelberg interferometry with a classical nanomechanical two-mode system

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Quantum Interfaces with Nano-opto-electro-mechanical devices: Applications and Fundamental Physics Ettore Majorana Foundation and Center for Scientific Culture, Erice, 2.8.2016

Classical coherence Interference of water waves

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1 µm



High stress = High Q Tensile stress increases the stored energy

Loss arises from anelasticity, i.e. a delay between internal strains and stresses: Assume complex Young's modulus $E = E_1 + i E_2$



see also: Gonzales & Saulson, J. Ac. Soc. Am. 96, 207 (1994) Unterreithmeier et al., Phys. Rev. Lett. 105, 027205 (2010) Yu et al., Phys. Rev. Lett. 108, 083603 (2012)

Ultra-high Q SiN resonators at 300 K Tensile stress of SiN film deposited on Si/SiO2 vs. fused silica wafer



Verbridge et al., J.Appl. Phys. 99, 124304 (2006) Faust, Krenn, Manus, Kotthaus, Weig, Nature Comm. 3, 728 (2012)





1. Nanomechanical SiN string resonators: Dielectric transduction and strong mode coupling



2. Landau-Zener dynamics: Coherent control of nanomechanical modes



3. Self-interference: Nanomechanical Stückelberg interferometry

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Operating high Q nanomechanics Using an electrical scheme... but without metallizing the resonator?

Metallizing the resonator enables electrical transduction



Operating high Q nanomechanics Using an electrical scheme... but without metallizing the resonator?

Metallizing the resonator enables electrical transduction but induces strong damping in metals at 300 K



> Avoid metallization-based transduction schemes for achieving highest Q

Seitner, Gajo, Weig, Appl. Phys. Lett. 105, 213101 (2014)

Dielectric gradient field transduction An integrated platform to control high Q nanomechanical resonators

Dielectric detection:

Heterodyne detection w/ 3.5 GHz microwave cavity



Dielectric mode coupling:

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Coupling spring κ induced by cross derivative of electric field



Faust, Phys. Rev. Lett 109, 037205 (2012)

Dielectric actuation:

Electrically induced gradient force

e.g. out of plane mode:

$$F_{z} = p_{y} \cdot \frac{\partial E_{z}}{\partial y}$$
$$\sim V_{DC}^{2} + 2V_{DC}V_{RF}$$



Unterreithmeier, Nature 458, 1001 (2009) see also: Schmid, APL 89, 163506 (2006)

Dielectric frequency tuning:

V_{DC}-controlled effective spring constant (local field gradient at string position)

e.g. out of plane mode:

$$\omega_{0} = \sqrt{\frac{\mathbf{k} + \mathbf{k}_{eff}}{\mathbf{m}}} \sim V_{DC}^{2}$$

w/ $\mathbf{k}_{eff} = -\frac{\partial F_{DC}}{\partial z}$



Rieger, Appl. Phys. Lett. 101, 103110 (2012)





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Tuning in- and out-of-plane flexural mode of a resonator in "elevated electrode" layout





Faust, Rieger, Seitner, Krenn, Manus, Kotthaus, Weig, Phys. Rev. Lett 109, 037205 (2012)

Tuning in- and out-of-plane flexural modes Avoided crossing reminiscent of strong coupling



Faust, Rieger, Seitner, Krenn, Manus, Kotthaus, Weig, Phys. Rev. Lett 109, 037205 (2012)

Time-resolved dynamics of coupled modes Measurement sequence

- State initialization at point | by constant drive
- 2. DC voltage ramp across coupling region
- Final state depends on ramp time τ: a diabatic / adiabatic transistion gets the system to point D / A
- 4. Measure oscillation energy at D and A (after delay δ)



A classical analogue of Landau-Zener physics Establishing time-domain control of nanoresonator state



Faust, Rieger, Seitner, Krenn, Manus, Kotthaus, Weig, Phys. Rev. Lett 109, 037205 (2012)

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A classical nanomechanical two-mode system and its two hybrid modes as basis states of a Bloch sphere

On resonance, the +45° and -45° mechanical hybrid modes form a two-mode system:



Faust, Rieger, Seitner, Kotthaus, Weig, Nature Physics 9, 485 (2013)

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Nanomechanical self-interference

induced by a double passage through the avoided crossing



Single passage through the avoided crossing:

 \rightarrow Landau-Zener dynamics

Double passage through the avoided crossing:

→ Stückelberg interference?

see: Stückelberg, Helv. Phys. Acta 5, 369 (1932)

 $\begin{array}{l} \rightarrow \text{Common approximation:} \\ \text{Adiabatic impulse model} \\ P_{1 \rightarrow 1} = 1 - 4P_{LZ} \left(1 - P_{LZ}\right) \sin^2 \left(\phi_{\text{s}} + \frac{\phi_{\text{dyn}}}{2} + \frac{\phi_{\text{geo}}}{2}\right) \\ \text{Review:} \\ \text{Shevchenko et al., Phys. Rep. 492, 1 (2010)} \end{array}$

Seitner et al., arXiv:1602.01034



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But would that also work by means of classical coherence? Remember the water waves...



Stückelberg interferometry with a classical two mode system Exact solution of the classical double passage problem



Two coupled harmonic oscillators:
$$\begin{split} m_{eff} \ddot{z} + m_{eff} \gamma_{1} \dot{z} + k_{1} \left(U_{DC} \right) z + \kappa \left(z - y \right) &= 0 \\ m_{eff} \ddot{y} + m_{eff} \gamma_{2} \dot{y} + k_{2} \left(U_{DC} \right) y + \kappa \left(y - z \right) &= 0 \end{split}$$

Normalized amplitudes: $y(t) = c_1(t) \cdot e^{i\omega_1 t}, z(t) = c_2(t) \cdot e^{i\omega_1 t}$ $2i\omega_1 \dot{c}_1 = \frac{\kappa}{m_{eff}} c_2$ $2i\omega_1 \dot{c}_2 + (\omega_2 - \omega_1)^2 c_2 = \frac{\kappa}{m_{eff}} c_1$

see: Novotny, Am. J. Phys. 78, 1199 (2010)

Time-dependend unitary transformation: © Hugo Ribeiro, McGill University

$$\mathbf{i} \begin{pmatrix} \dot{\mathbf{C}}_1 \\ \dot{\mathbf{C}}_2 \end{pmatrix} = \begin{pmatrix} \frac{\alpha t}{2} & \frac{\lambda}{2} \\ \frac{\lambda}{2} & -\frac{\alpha t}{2} \end{pmatrix} \begin{pmatrix} \mathbf{C}_1 \\ \mathbf{C}_2 \end{pmatrix}$$

EOM of two coupled harmonic oscillators looks like Schrödinger equation of the Landau-Zener problem!

see: Vitanov & Garraway, Phys. Rev. A 53, 4288 (1996)

Return probability $P_{1\rightarrow 1}$ is identical in the classical and the quantum case.

 $P_{1 \to 1} = \left| \phi_{11} \left(t_{p}, t_{i} \right) \phi_{11}^{*} \left(t, -t_{p} \right) + \phi_{12}^{*} \left(t_{p}, t_{i} \right) \phi_{12}^{*} \left(t, -t_{p} \right) \right|^{2}$

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Classical Stückelberg interference at 10 K and analytical model by Hugo Ribeiro



theory (no free parameters)

Seitner, Ribeiro, Kölbl, Faust, Kotthaus, Weig, arXiv:1602.01034

Classical Stückelberg interference at 300 K indicating millisecond coherence time



Seitner, Ribeiro, Kölbl, Faust, Kotthaus, Weig, arXiv:1602.01034

Classical Stückelberg interference at room temperature indicating millisecond coherence time



Discrepancies:

- Temperature drifts:
 ±2 K/h and
 corresponding shift
 of eigenfrequencies
 by 500 Hz/K, i.e.
 40 linewidths affects
 all parameters of the
 system
- Feedback loop ensures initialization at fixed frequecy
- Other parameters not precisely known (e.g. U_{ac})

Classical to quantum transition

Would this also work with a mechanical resonator in the quantum regime?



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SUMMARY



- High stress SiN nanomechanical string resonators:
 - Q > 300,000 at T = 300 K
 - Excellent control via dielectric transduction
 - Versatile toolbox for nanomechanics



- Strong coupling between in- and out-of-plane mode:
 - \blacktriangleright g/ Γ ~ 10² mediated by inhomogeneous electric field
 - Classical Landau-Zener dynamics
 - Coherent control of nanomechanical motion



- Classical Stückelberg interferometry:
 - Same return probability as in quantum mechanical case
 - Observation of finite time effects
 - Nanomechanical interferometer

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